

Lecture 10:

Numerical Spectral method

Since $\{ \overrightarrow{e^{ikx}} \}_{k=0}^{N-1}$ is a basis. We can write:

$$\vec{u} = \sum_{k=0}^{N-1} \hat{u}_k \overrightarrow{e^{ikx}} \quad \text{and} \quad \vec{f} = \sum_{k=0}^{N-1} \hat{f}_k \overrightarrow{e^{ikx}}$$

In other words, for each j , $f_j = f(x_j) = \sum_{k=0}^{N-1} \hat{f}_k (e^{ikx})_j$

$$= \sum_{k=0}^{N-1} \hat{f}_k e^{ikx_j} = \sum_{k=0}^{N-1} \hat{f}_k e^{i k j \left(\frac{2\pi}{N}\right)}$$

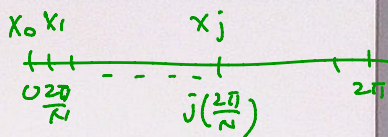
$\therefore \hat{f}_k$ can be determined by DFT.

To solve $\frac{d^2 u}{dx^2} = f$, we approximate it by

$$\hat{\vec{f}} = \begin{pmatrix} \hat{f}_0 \\ \hat{f}_1 \\ \vdots \\ \hat{f}_{N-1} \end{pmatrix} = \frac{A_w}{N} \vec{f}$$

$\hat{\vec{D}} \vec{u} = \hat{\vec{f}}$

$$\mathbb{C}^N \ni \overrightarrow{e^{ikx}} = \begin{pmatrix} e^{ikx_0} \\ e^{ikx_1} \\ \vdots \\ e^{ikx_{N-1}} \end{pmatrix}$$



Now, $\tilde{D} \vec{u} = \vec{f}$ becomes:

$$\tilde{D} \left(\sum_{k=0}^{N-1} \hat{u}_k \vec{e}^{ikx} \right) = \sum_{k=0}^{N-1} \hat{f}_k \vec{e}^{ikx}$$

$$\Leftrightarrow \sum_{k=0}^{N-1} \hat{u}_k \tilde{D} \vec{e}^{ikx} = \sum_{k=0}^{N-1} \hat{f}_k \vec{e}^{ikx}$$

\parallel
 $(-\lambda_k^2) \vec{e}^{ikx}$

$$\Leftrightarrow \sum_{k=0}^{N-1} \hat{u}_k (-\lambda_k^2) \vec{e}^{ikx} = \sum_{k=0}^{N-1} \hat{f}_k \vec{e}^{ikx}$$

Comparing coefficients, we get

$$\underbrace{-\lambda_k^2}_{\text{known}} \underbrace{\hat{u}_k}_{\text{unknown}} = \underbrace{\hat{f}_k}_{\text{known}} \quad \text{for } k=0, 1, 2, \dots, N-1$$

(algebraic equation)

For $k=1, 2, \dots, N-1$, we have: $\hat{u}_k = \hat{f}_k / (-\lambda_k^2)$.

For $k=0$, $\lambda_k=0$!!

We consider a special solution such that:

$$\hat{u}_0 = \underbrace{u_0 + u_1 + \dots + u_{N-1}}_N = 0 \quad (\text{Zero-mean solution})$$

Then, we can set $\hat{u}_0 = 0$

$$\text{Note that } \hat{f}_0 = -\lambda_0^2 \hat{u}_0 = 0 \Rightarrow \frac{f_0 + f_1 + \dots + f_{N-1}}{N} = 0$$

This is consistent with the compatible condition of f :

$$\int_0^{2\pi} f(x) dx = \int_0^{2\pi} u''(x) dx = u'(x) \Big|_0^{2\pi} = 0 \quad (\text{Periodic})$$

\Downarrow
 \hat{f}_0

Once \hat{u}_k is obtained for $k=0, 1, 2, \dots, N-1$, \vec{u} can be obtained

$$\text{by: } \vec{u} = \begin{pmatrix} u(x_0) \\ u(x_1) \\ \vdots \\ u(x_{N-1}) \end{pmatrix} = \sum_{k=0}^{N-1} \hat{u}_k \overrightarrow{e^{ikx}} \quad (\text{inverse DFT})$$

$$\text{or } \vec{u} = A\omega \begin{pmatrix} \hat{u}_0 \\ \hat{u}_1 \\ \vdots \\ \hat{u}_{N-1} \end{pmatrix} \quad \text{where } A\omega = \begin{pmatrix} 1 & 1 & \dots & 1 \\ 1 & \omega & & \omega^{N-1} \\ \vdots & \omega^2 & & \vdots \\ \vdots & \vdots & & \vdots \\ 1 & \omega^{N-1} & & \omega^{(N-1)^2} \end{pmatrix}$$

$\omega = e^{i\frac{2\pi}{N}}$

(Matrix multiplication)

Remark: For any other solution \vec{u}_* , $\vec{u}_* = \vec{u} + c \begin{pmatrix} 1 \\ \vdots \\ 1 \end{pmatrix}$ for some constant c .

c is determined by other condition (such as boundary condition)

Example: Consider: $a \frac{d^2 u}{dx^2} + b \frac{du}{dx} = f$ for $x \in [0, 2\pi]$.

This time, we approximate $\frac{d^2 u}{dx^2}$ by =

$$(*) \quad \frac{d^2 u}{dx^2}(x_j) \approx \frac{u_{j-2} - 2u_j + u_{j+2}}{4h^2} \quad \text{for } j = 0, 1, 2, \dots, N-1$$

Again, we assume $u_{-1} = u_{N-1}$, $u_1 = u_{N+1}$, $u_{-2} = u_{N-2}$, ..., etc

Motivation: ① $u(x_j + 2h) \approx u(x_j) + 2h u'(x_j) + 2h^2 u''(x_j)$

② $u(x_j - 2h) \approx u(x_j) - 2h u'(x_j) + 2h^2 u''(x_j)$

$$① + ② : \quad u(x_{j+2}) + u(x_{j-2}) - 2u(x_j) = 4h^2 u''(x_j)$$

This time, we approximate $\frac{du}{dx}$ as =

$$(**) \quad \frac{du}{dx}(x_j) \approx \frac{u_{j+1} - u_{j-1}}{2h}$$

Claim: $\overrightarrow{e^{ikx}}$ is an eigenvector of \tilde{D} and D .

Proof: $(D \overrightarrow{e^{ikx}})_j = \frac{e^{ikx_{j-2}} - 2e^{ikx_j} + e^{ikx_{j+2}}}{4h^2}$

$$= \frac{e^{ikx_j} (e^{-2ikh} - 2 + e^{2ikh})}{4h^2}$$

$$= \underbrace{\left(\frac{-\sin^2(kh)}{h^2} \right)}_{\tilde{\lambda}_k^2} e^{ikx_j} \quad \therefore D \overrightarrow{e^{ikx}} = \tilde{\lambda}_k \overrightarrow{e^{ikx}}$$

Also, $\tilde{D} \overrightarrow{e^{ikx}} = \frac{e^{ikx_{j+1}} - e^{ikx_{j-1}}}{2h} = e^{ikx_j} (e^{ikh} - e^{-ikh})$

$$\therefore \tilde{D} \overrightarrow{e^{ikx}} = \tilde{\lambda}_k \overrightarrow{e^{ikx}} = \underbrace{\left(\frac{i \sin(kh)}{h} \right)}_{\tilde{\lambda}_k} e^{ikx_j}$$

Now, $a \frac{d^2 u}{dx^2} + b \frac{du}{dx} = f$ can be discretized as:

$$a \tilde{D}^2 \vec{u} + b \tilde{D} \vec{u} = \vec{f} \quad \text{subject to the periodic condition.}$$

Recall: $\{e^{ikx}\}_{k=0}^{N-1}$ is a basis for \mathbb{C}^N

Again, let $\vec{u} = \sum_{k=0}^{N-1} \hat{u}_k e^{ikx}$ and $\vec{f} = \sum_{k=0}^{N-1} \hat{f}_k e^{ikx}$

$$\text{Then: } a \tilde{D}^2 \left(\sum_{k=0}^{N-1} \hat{u}_k e^{ikx} \right) + b \tilde{D} \left(\sum_{k=0}^{N-1} \hat{u}_k e^{ikx} \right) = \sum_{k=0}^{N-1} \hat{f}_k e^{ikx}$$

$$\Leftrightarrow \sum_{k=0}^{N-1} (a \tilde{\lambda}_k^2 + b \tilde{\lambda}_k) \hat{u}_k e^{ikx} = \sum_{k=0}^{N-1} \hat{f}_k e^{ikx}$$

Comparing coefficients, we get $(a \tilde{\lambda}_k^2 + b \tilde{\lambda}_k) \hat{u}_k = \hat{f}_k$ for $k=0, \dots, N-1$
(algebraic equation)

For $k=0$ and $\frac{N}{2}$, $\tilde{\lambda}_k = 0$. We set $\hat{u}_0 = 0 = \hat{u}_{\frac{N}{2}}$.

In general, we set:
$$\hat{u}_k = \begin{cases} \hat{f}_k / (a \tilde{\lambda}_k^2 + b \tilde{\lambda}_k) & \text{for } k \neq 0, \frac{N}{2} \\ 0 & \text{for } k=0 \text{ or } k = \frac{N}{2} \end{cases}$$

\vec{u} can be obtained by inverse DFT =

$$\begin{pmatrix} u_0 \\ u_1 \\ \vdots \\ u_{N-1} \end{pmatrix} = A \omega \begin{pmatrix} \hat{u}_0 \\ \hat{u}_1 \\ \vdots \\ \hat{u}_{N-1} \end{pmatrix}$$

Question: How about the general solution?

Answer: Examine $N(a \tilde{D}^2 + b \tilde{D})$, the null space.

Claim: $N(a\tilde{D}^2 + b\tilde{D}) = \text{span} \{ \overline{e^{i(0)x}}, \overline{e^{i(\frac{N}{2})x}} \}$

$$= \text{span} \left\{ \begin{pmatrix} | \\ | \\ \vdots \\ | \\ | \end{pmatrix}, \begin{pmatrix} e^{i(\frac{N}{2})x_1} \\ \vdots \\ e^{i(\frac{N}{2})x_{N-1}} \end{pmatrix} \right\}$$

Proof: $a\tilde{D}^2 + b\tilde{D}$ has two eigenvectors whose eigenvalue is 0.

These eigenvectors are $\overline{e^{i(0)x}}$ and $\overline{e^{i(\frac{N}{2})x}}$.

$N(a\tilde{D}^2 + b\tilde{D}) =$ eigenspace of eigenvalue = 0

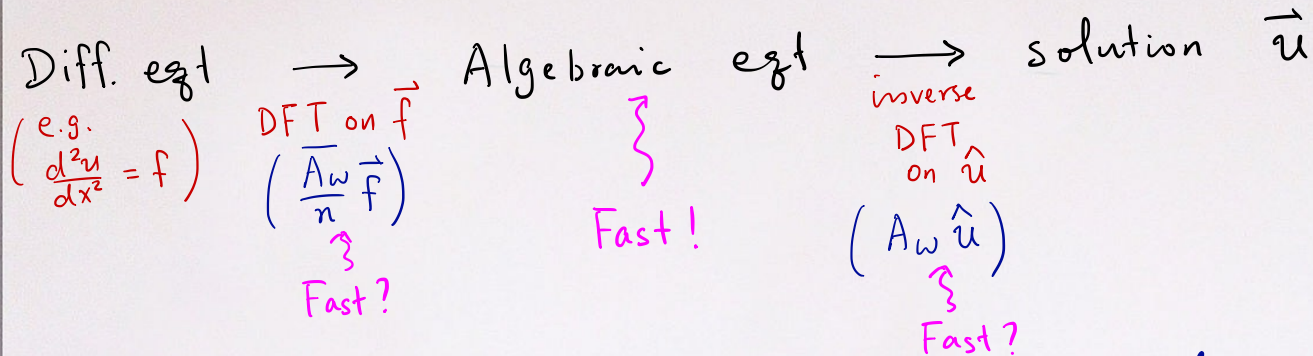
$$= \text{span} \{ \overline{e^{i(0)x}}, \overline{e^{i(\frac{N}{2})x}} \}$$

Thus, general sol: $\vec{u}^x = \vec{u} + c_1 \overline{e^{i(0)x}} + c_2 \overline{e^{i(\frac{N}{2})x}}$

for some c_1 and c_2 .

c_1 and c_2 can be determined by certain conditions (such as boundary conditions)

Main idea of numerical spectral method



Remark: To develop an efficient numerical spectral method, we need to compute $A_w \hat{u}$ and $\frac{\overline{A_w} \vec{f}}{n}$ fast.

- Computational cost for $A_w \hat{u}$ is $\mathcal{O}(n^2)$.
($n \times n$)

Goal: Reduce the computational cost to $\mathcal{O}(n \log n)$

e.g. $n = 2^{10}$, $n^2 = 2^{20}$, $n \log n = 10 \cdot 2^{10} < 2^{14}$. $\therefore 2^6 = 64$ times faster!